

1. Subject: Transonic Potential Equations

The Transonic Potential Equations or the Linear Cauchy-Riemann Equations are formed from the coupling of the steady compressible continuity equation of fluid dynamics

$$\partial_x \rho u + \partial_y \rho v = 0 \quad (1)$$

and the vorticity definition

$$-\partial_x v + \partial_y u = 0 \quad (2)$$

with vorticity  $\omega = 0$  the irrotational potential assumption. Here  $\rho$  is density,  $u, v$  the Cartesian velocity components and the isentropic assumption leads us to

$$\rho = \left(1 - \frac{\gamma - 1}{2} (u^2 + v^2 - 1)\right)^{\frac{1}{\gamma - 1}} \quad (3)$$

with  $\gamma = 1.4$  the ratio of specific heats. Pressure is defined as  $p = \rho^\gamma$ .

The linear Cauchy-Riemann equations are recovered above and below by setting  $\rho = 1$  instead of Eq 3.

Combining Eq 1 and Eq 2 in vector form we have

$$\partial_x \mathbf{f}(\mathbf{q}) + \partial_y \mathbf{g}(\mathbf{q}) = 0 \quad (4)$$

where

$$\mathbf{q} = \begin{pmatrix} u \\ v \end{pmatrix}, \quad \mathbf{f} = \begin{pmatrix} -\rho u \\ v \end{pmatrix}, \quad \mathbf{g} = \begin{pmatrix} -\rho v \\ -u \end{pmatrix} \quad (5)$$

**Note:** One approach to solving these equations is to cast them as a hyperbolic system where we solve

$$\partial_t \mathbf{q} + \partial_x \mathbf{f} + \partial_y \mathbf{g} = 0 \quad (6)$$

(a) For Eq 6

- i. Find the flux Jacobians of  $\mathbf{f}$  and  $\mathbf{g}$ .
- ii. Determine the eigenvalues and conditions under which the system is hyperbolic. (Hint: A system is hyperbolic if the eigenvalues of it's flux Jacobians are real.)

**ANSWER:** Define  $\frac{\partial \mathbf{f}}{\partial \mathbf{q}} = A$  and  $\frac{\partial \mathbf{g}}{\partial \mathbf{q}} = B$

The Jacobian matrices  $A$  and  $B$  are

$$A = \begin{bmatrix} -\rho + \rho^{2-\gamma} u^2 & \rho^{2-\gamma} uv \\ 0 & 1 \end{bmatrix}$$

and

$$B = \begin{bmatrix} \rho^{2-\gamma} uv & -\rho + \rho^{2-\gamma} v^2 \\ -1 & 0 \end{bmatrix}$$

The eigenvalues of the Jacobian matrices are interesting and shed some light on the form of the scheme's stability and formulation. The eigenvalues of  $A$  and  $B$  are

$$\lambda(A) = -\rho + \rho^{2-\gamma} u^2, \quad 1$$

$$\lambda(B) = \frac{\rho^{2-\gamma}uv \pm \sqrt{\rho^{4-2\gamma}u^2v^2 + 4\rho - 4\rho^{2-\gamma}v^2}}{2}$$

For the system to be hyperbolic the  $\lambda$ 's must be real. Since  $\rho > 0$ , this implies that

$$\rho^{4-2\gamma}u^2v^2 + 4\rho \geq 4\rho^{2-\gamma}v^2$$

**Discussion:** We can examine this condition in light of assumptions on the values of  $u$  and  $v$ . If  $v = 0$  the condition is always satisfied, if  $u = 0$  then  $\rho^{1-\gamma}v^2 < 1$  which for  $0.2 \leq \rho \leq 2$  (quite a large range of  $\rho$ ) implies that  $0 < |v| < 1.9$ , a reasonable range. What has not been discussed here, is that the full potential equations are limited to weak shock waves, which translates in this case to  $|u|, |v| \leq 2.0$ , as a general rule.

(b) The fluxes of the Euler equations are homogeneous of degree 1.

i. Are the above fluxes  $\mathbf{f}$  and  $\mathbf{g}$  homogeneous of degree 1?, degree  $n$ ?

ii. If we replace Eq 3 with  $\rho = 1$ , what can be said about the properties of the system?

The fluxes in Eq. 4 are not homogeneous because of the nonlinear nature of Eq. 3.

$$\mathbf{f}(\alpha \mathbf{q}) = \left( - \left( 1 - \frac{\gamma-1}{2} \alpha^2 (u^2 + v^2 - 1) \right)^{\frac{1}{\gamma-1}} u \right) \neq \alpha \mathbf{f}(\mathbf{q})$$

This does not satisfy the homogeneity property, for any  $n$ .

For  $\rho = 1$ ,

$$\mathbf{f}(\alpha \mathbf{q}) = \begin{pmatrix} -\alpha u \\ \alpha v \end{pmatrix} = \alpha \begin{pmatrix} -u \\ v \end{pmatrix} = \alpha \mathbf{f}(\mathbf{q})$$

and

$$\mathbf{g}(\alpha \mathbf{q}) = \begin{pmatrix} -\alpha v \\ -\alpha u \end{pmatrix} = \alpha \begin{pmatrix} -v \\ -u \end{pmatrix} = \alpha \mathbf{g}(\mathbf{q})$$

The system is linear and homogeneous of degree 1.

(a) Subject: Splitting/Factorization

i. *Space vector definition* The space vector for the natural ordering is

$$u^{(n)} = [u_1 \ u_2 \ u_3 \ u_4 \ u_5 \ u_6 \ u_7 \ u_8]^T$$

ii. The space vector for an odd-even ordering is

$$u^{(*)} = [u_1 \ u_3 \ u_5 \ u_7 \ u_2 \ u_4 \ u_6 \ u_8]^T$$

iii.  $P_{n*}$  is the permutation matrix that permutes the odd-even ordering to the natural ordering, and  $P_{*n}$  is the permutation matrix that permutes the natural ordering to the odd-even ordering. Simply put

$$\begin{aligned} u^{(n)} &= P_{n*} u^{(*)} \\ u^{(*)} &= P_{*n} u^{(n)} \end{aligned}$$

The permutation matrices are easily deduced

$$P_{n*} = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix} \quad \text{and} \quad P_{*n} = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix}$$

Note that these permutation matrices are the transpose of each other and the inverse of each other.

iv.  $A^{(n)}$  is a simple tridiagonal matrix

$$A^{(n)} = \frac{\nu}{\Delta x^2} \begin{bmatrix} -2 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & -2 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & -2 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -2 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & -2 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & -2 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & -2 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & -2 \end{bmatrix}$$

Premultiplying the generic ODE in naturally ordered form by  $P_{*n}$  gives

$$P_{*n} \frac{du^{(n)}}{dt} = P_{*n} A^{(n)} u^{(n)} + P_{*n} f$$

Since the permutation matrices are not functions of time,  $P_{n*}$  can be brought inside the time derivative on the *LHS* of the above equation to give

$$\frac{dP_{*n} u^{(n)}}{dt} = P_{*n} A^{(n)} u^{(n)} + P_{*n} f$$

Noting that  $P_{n*} P_{*n} = I$ , and that multiplying a term by the identity matrix leaves the term unchanged, the above equation can be written as

$$\frac{dP_{*n} u^{(n)}}{dt} = P_{*n} A^{(n)} P_{n*} P_{*n} u^{(n)} + P_{*n} f$$

Or more succinctly as

$$\frac{du^{(*)}}{dt} = P_{*n} A^{(n)} P_{n*} u^{(*)} + P_{*n} f$$

Comparing with the odd-even version of the ODE

$$\frac{du^{(*)}}{dt} = A^{(*)} u^{(*)} + g$$

it readily follows that  $A^{(*)} = P_{*n} A^{(n)} P_{n*}$  and  $g = P_{*n} f = 0$ . So the odd-even ordered spatial matrix operator is

$$A^{(*)} = \frac{\nu}{\Delta x^2} \begin{bmatrix} -2 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & -2 & 0 & 0 & 1 & 1 & 0 & 0 \\ 0 & 0 & -2 & 0 & 0 & 1 & 1 & 0 \\ 0 & 0 & 0 & -2 & 0 & 0 & 1 & 1 \\ 1 & 1 & 0 & 0 & -2 & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 & 0 & -2 & 0 & 0 \\ 0 & 0 & 1 & 1 & 0 & 0 & -2 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & -2 \end{bmatrix}$$

v. In an implicit Euler time differencing formulation of the generic ODE the time derivative is first order accurate, and the spatial term,  $Au$ , is evaluated at time  $n + 1$ .

$$\frac{du}{dt} = \frac{1}{h} (u_{n+1} - u_n) + O(h) = Au_{n+1}$$

or, equivalently,

$$[I - hA]u_{n+1} = u_n + O(h^2)$$

Subtracting  $[I - hA]u_n$  from both sides, and noting that  $\Delta u_n \equiv u_{n+1} - u_n$ , gives

$$[I - hA]\Delta u_n = hAu_n + O(h^2)$$

This is the delta form of the implicit algorithm. For the naturally ordered form the implicit matrix operator is a simple tridiagonal matrix; for the odd-even form it's a banded matrix with a bandwidth of nine.

(b) *System definition*

- i. Define the odd and even matrix operators as follows

$$A^{(o)} = \frac{\nu}{\Delta x^2} \begin{bmatrix} -2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -2 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -2 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -2 & 0 & 0 & 0 & 0 \\ 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \end{bmatrix}$$

$$A^{(e)} = \frac{\nu}{\Delta x^2} \begin{bmatrix} 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 0 & -2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -2 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -2 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -2 \end{bmatrix}$$

The resulting ODE is

$$\frac{du^{(*)}}{dt} = [A^{(o)} + A^{(e)}] u^{(*)}$$

- ii. Define the diagonal matrix,  $D$ , and the upper matrix,  $U$ , as

$$D = \frac{\nu}{\Delta x^2} \begin{bmatrix} -2 & 0 & 0 & 0 \\ 0 & -2 & 0 & 0 \\ 0 & 0 & -2 & 0 \\ 0 & 0 & 0 & -2 \end{bmatrix} \quad \text{and} \quad U = \frac{\nu}{\Delta x^2} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

It readily follows that

$$A^{(o)} = \left[ \begin{array}{c|c} D & 0_4 \\ \hline U & 0_4 \end{array} \right] \quad \text{and} \quad A^{(e)} = \left[ \begin{array}{c|c} 0_4 & U^T \\ \hline 0_4 & D \end{array} \right]$$

- iii. Results from problem 1a 2a are easily adapted to give

$$[I - hA^{(o)} - hA^{(e)}] u_{n+1}^{(*)} = u_n^{(*)} + O(h^2)$$

and in delta form

$$[I - hA^{(o)} - hA^{(e)}] \Delta u_n^{(*)} = [hA^{(o)} + hA^{(e)}] u_n^{(*)} + O(h^2)$$

and in factored delta form

$$[I - hA^{(o)}] [I - hA^{(e)}] \Delta u_n^{(*)} = [hA^{(o)} + hA^{(e)}] u_n^{(*)} + O(h^2)$$

Since  $h^2 A^{(o)} A^{(e)} \Delta u_n^{(*)}$  is  $O(h^3)$ , the error term in the factored form remains  $O(h^2)$ .

iv. The odd and even matrix factors are

$$\left[ I - hA^{(o)} \right] = \left[ \begin{array}{c|c} I_4 - hD & 0_4 \\ \hline -hU & I_4 \end{array} \right] \quad \text{and} \quad \left[ I - hA^{(e)} \right] = \left[ \begin{array}{c|c} I_4 & -hU^T \\ \hline 0_4 & I_4 - hD \end{array} \right]$$

So clearly the odd matrix factor is a lower triangular matrix, and the even matrix factor is an upper triangular matrix. The factoring produced an  $LU$  decomposition of the implicit matrix operator for the odd-even ordering of the system without degrading the order of accuracy. The system may now be solved by simple forward and backward substitutions: Let

$$\tilde{u}_{n+1}^{(*)} = U u_{n+1}^{(*)} = \left[ I - hA^{(e)} \right] u_{n+1}^{(*)}$$

then

$$L \tilde{u}_{n+1}^{(*)} = \left[ I - hA^{(o)} \right] \tilde{u}_{n+1}^{(*)} = \left[ hA^{(o)} + hA^{(e)} \right] u_n^{(*)}$$

$\tilde{u}_{n+1}^{(*)}$  may be solved for by forward substitution, and then  $u_{n+1}^{(*)}$  can be solved for by backward substitution. This solution process offers a considerable savings, in terms of multiplication and division operations count, over the direct inversion of the resulting system from problem 1(a)v. Direct inversion using Gaussian elimination for an  $m \times m$  nonsymmetric fully populated matrix requires  $m(2m-1)(2m+1)/3$  multiplication and division operations, whereas forward and backward substitutions each require  $m(m+1)/2$  multiplication and division operations.